

Floquet engineering for quantum simulation

MAX PLANCK INSTITUTE FOR THE PHYSICS OF COMPLEX SYSTEMS





lab frame

rotating frame





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MPI-PKS (Dresden)



Quantum technologies

Quantum Communication



image: PRL 123 100506 (2019)

process info with unprecedented security

Quantum Computing



image: IBM

speed up essential algorithms

Quantum Sensing & Metrology



image: ETH Zurich

measure weakest of fields

Quantum Simulation



 understand properties of quantum matter, complex molecules, drug discovery



superconducting qubits



image: IBM

neutral atoms



trapped ions



photons

image: ParityQC

image: Wikipedia





image: arXiv:2404.05620





Quantum Simulators

quantum simulator







Richard P Feynman

"Nature isn't classical, dammit, and if you want to make a simulation of nature, you'd better make it quantum mechanical, and by golly it's a wonderful problem, because it doesn't look so easy." (1982)

- use one quantum system to emulate the behavior of another
- restrictions: not all quantum systems can be simulated

Q: how can we expand the range of systems we can simulate?

Marín Bukov

Gross & Bloch, Science 357, 6355 (2017)



Periodically driven systems



video: YouTube (bluedwarf1127)



video: YouTube (Harvard Nat Sci)

High-frequency periodic drives can change drastically the fundamental properties of physical systems







- classical systems: fictitious forces
- quantum systems
- Periodically driven quantum systems
 - Floquet theorem
 - Floquet engineering
- Examples
 - spin-1 particle in a circularly polarized drive
 - quantum Kapitza oscillator
 - artificial gauge fields











- Rotating reference frames
 - classical systems: fictitious forces



lab frame

rotating frame



- rotating reference frame
 - not inertial
 - fictitious forces

Merry-go-round







- rotating reference frame
 - not inertial

fictitious forces

e.g.,
$$P(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0 \\ \sin \omega t & \cos \omega t & 0 \\ 0 & 0 & 1 \end{pmatrix}$$



transformation between lab and rotating frames



rotation matrix: P(t)

lab frame

rot frame







 $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$





 $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$



 rotating reference frame $\hat{z} = \hat{z}_{rot}$ e.g., $P(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0 \\ \sin \omega t & \cos \omega t & 0 \\ 0 & 0 & 1 \end{pmatrix}$ not inertial $\hat{y}_{rot}(t)$ fictitious forces transformation between lab and rotating frames $P^{-1}(t) = P^{\dagger}(t)$ lab frame rot frame rotation matrix: P(t)position: $\vec{r}_{lab}(t)$ $\vec{r}_{\rm rot}(t)$ $\dot{\vec{r}}_{\rm rot}(t)$ velocity: $\dot{\vec{r}}_{lab}(t)$ $\ddot{\vec{r}}_{\rm rot}(t)$ acceleration: $\vec{r}_{lab}(t)$ $\overrightarrow{F}_{\rm rot}$ force: \overrightarrow{F}_{lab} $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$ $\vec{F}_{lab}(t) = P(t) \vec{F}_{rot}(t)$



• time-dependent rotation matrix P(t)

$$f(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0\\ \sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{z} = \hat{z}_{rot}$$

$$\hat{y}_{rot}(t)$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{y}_{rot}(t)$$



- time-dependent rotation matrix $P(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0\\ \sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix} \hat{x}$
 - inverse transformation

$$P^{-1}(t) = P^{\dagger}(t) = P(t; -\omega) = \begin{pmatrix} \cos \omega t & +\sin \omega t & 0\\ -\sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

$$\hat{z} = \hat{z}_{\text{rot}}$$

$$\hat{y}_{\text{rot}}(t)$$

$$\hat{y}_{\text{rot}}(t)$$



• time-dependent rotation matrix

$$P(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0\\ \sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

inverse transformation

$$P^{-1}(t) = P^{\dagger}(t) = P(t; -\omega) = \begin{pmatrix} \cos \omega t & +\sin \omega t & 0\\ -\sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

element-wise derivatives

$$\dot{P}(t) = \omega \begin{pmatrix} -\sin\omega t & -\cos\omega t & 0\\ \cos\omega t & -\sin\omega t & 0\\ 0 & 0 & 0 \end{pmatrix}$$





• time-dependent rotation matrix

$$P(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0\\ \sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

inverse transformation

$$P^{-1}(t) = P^{\dagger}(t) = P(t; -\omega) = \begin{pmatrix} \cos \omega t & +\sin \omega t & 0\\ -\sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

► element-wise derivatives
$$\dot{P}(t) = \omega \begin{pmatrix} -\sin \omega t & -\cos \omega t & 0 \\ \cos \omega t & -\sin \omega t & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$P^{\dagger}(t)\dot{P}(t) = \begin{pmatrix} \cos\omega t & +\sin\omega t & 0\\ -\sin\omega t & \cos\omega t & 0\\ 0 & 0 & 1 \end{pmatrix} \omega \begin{pmatrix} -\sin\omega t & -\cos\omega t & 0\\ \cos\omega t & -\sin\omega t & 0\\ 0 & 0 & 0 \end{pmatrix} = \omega \begin{pmatrix} 0 & -1 & 0\\ 1 & 0 & 0\\ 0 & 0 & 0 \end{pmatrix} = \omega \hat{z} \times \begin{pmatrix} 0 & -1 & 0\\ 1 & 0 & 0\\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} a\\ b\\ c \end{pmatrix} = \begin{pmatrix} -b\\ a\\ 0 \end{pmatrix} = \begin{pmatrix} 0\\ 0\\ 1 \end{pmatrix} \times \begin{pmatrix} a\\ b\\ c \end{pmatrix}$$

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 $\hat{z} = \hat{z}_{\rm rot}$

 $\hat{x}_{\rm rot}(t)$

â



time-dependent rotation matrix

$$P(t) = \begin{pmatrix} \cos \omega t & -\sin \omega t & 0\\ \sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

$$P^{-1}(t) = P^{\dagger}(t) = P(t; -\omega) = \begin{pmatrix} \cos \omega t & +\sin \omega t & 0\\ -\sin \omega t & \cos \omega t & 0\\ 0 & 0 & 1 \end{pmatrix}$$

► element-wise derivatives
$$\dot{P}(t) = \omega \begin{pmatrix} -\sin \omega t & -\cos \omega t & 0 \\ \cos \omega t & -\sin \omega t & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$P^{\dagger}(t)\dot{P}(t) = \begin{pmatrix} \cos \omega t & +\sin \omega t & 0 \\ -\sin \omega t & \cos \omega t & 0 \\ 0 & 0 & 1 \end{pmatrix} \omega \begin{pmatrix} -\sin \omega t & -\cos \omega t & 0 \\ \cos \omega t & -\sin \omega t & 0 \\ 0 & 0 & 0 \end{pmatrix} = \omega \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} = \omega \hat{z} \times$$

• general time-dependent rotation axis $\vec{\omega}(t)$ $\begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} a \\ b \\ c \end{pmatrix} = \begin{pmatrix} -b \\ a \\ 0 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} \times \begin{pmatrix} a \\ b \\ c \end{pmatrix}$

$$P^{\dagger}(t)\dot{P}(t) = \vec{\omega}(t) \times$$

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 $\hat{z} = \hat{z}_{\rm rot}$

â





Newton's law

• lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t) / P^{\dagger}(t) \cdot$$

 $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$ $\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$



Newton's law

► lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t) / P^{\dagger}(t) \cdot$$

• rot frame:
$$mP^{\dagger}(t)\frac{d^2}{dt^2}\left[P(t)\underline{P^{\dagger}(t)\vec{r}_{lab}(t)}\right] = P^{\dagger}(t)\vec{F}_{lab}(t)$$



$\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$ $\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$



Newton's law

► lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t) / P^{\dagger}(t) \cdot$$

• rot frame:
$$mP^{\dagger}(t)\frac{d^2}{dt^2}\left[P(t)\underline{P^{\dagger}(t)}\vec{r}_{lab}(t)\right] = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$mP^{\dagger}(t)\frac{\mathrm{d}}{\mathrm{d}t}\left(P(t)P^{\dagger}(t)\frac{\mathrm{d}}{\mathrm{d}t}\left[P(t)\vec{r}_{\mathrm{rot}}(t)\right]\right) = P^{\dagger}(t)\vec{F}_{\mathrm{lab}}(t)$$

$$\hat{z} = \hat{z}_{\text{rot}}$$

$$\hat{y}_{\text{rot}}(t)$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}_{\text{rot}}(t)$$

 $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$ $\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$



Newton's law

► lab frame:
$$m \frac{\mathrm{d}^2}{\mathrm{d}t^2} \vec{r}_{\mathrm{lab}}(t) = \vec{F}_{\mathrm{lab}}(t)$$
 / $P^{\dagger}(t)$.

• rot frame:
$$mP^{\dagger}(t)\frac{d^2}{dt^2}\left[P(t)P^{\dagger}(t)\vec{r}_{lab}(t)\right] = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$\hat{z} = \hat{z}_{rot}$$

$$\hat{y}_{rot}(t)$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}_{rot}(t)$$

$$\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$$
$$\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$$

$$mP^{\dagger}(t)\frac{d}{dt}\left(P(t)P^{\dagger}(t)\frac{d}{dt}\left[P(t)\vec{r}_{rot}(t)\right]\right) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$m\left(P^{\dagger}(t)\frac{d}{dt}P(t)\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$P^{\dagger}(t)\frac{d}{dt}P(t)\vec{f}(t) = P^{\dagger}(t)\vec{P}(t)\vec{f}(t) + \frac{d}{dt}\vec{f}(t)$$

$$= \left(P^{\dagger}(t)\dot{P}(t) + \frac{d}{dt}\right)\vec{f}(t)$$



• Newton's law

► lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t) / P^{\dagger}(t) \cdot$$

• rot frame:
$$mP^{\dagger}(t)\frac{\mathrm{d}^2}{\mathrm{d}t^2}\left[P(t)P^{\dagger}(t)\vec{r}_{\mathrm{lab}}(t)\right] = P^{\dagger}(t)\vec{F}_{\mathrm{lab}}(t)$$

$$mP^{\dagger}(t)\frac{d}{dt}\left(P(t)P^{\dagger}(t)\frac{d}{dt}\left[P(t)\vec{r}_{rot}(t)\right]\right) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$m\left(P^{\dagger}(t)\frac{d}{dt}P(t)\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$m\left(\underline{P^{\dagger}(t)\dot{P}(t)} + \frac{d}{dt}\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$\hat{z} = \hat{z}_{rot}$$

$$\hat{y}_{rot}(t)$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}_{rot}(t)$$

$$\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$$
$$\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$$

$${}^{\dagger}(t)\frac{\mathrm{d}}{\mathrm{d}t}P(t)\vec{f}(t) = P^{\dagger}(t)\dot{P}(t)\vec{f}(t) + \frac{\mathrm{d}}{\mathrm{d}t}\vec{f}(t)$$
$$= \left(P^{\dagger}(t)\dot{P}(t) + \frac{\mathrm{d}}{\mathrm{d}t}\right)\vec{f}(t)$$

$$P^{\dagger}(t)\dot{P}(t) = \overrightarrow{\omega}(t) \times$$

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Newton's law

• lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t) / P^{\dagger}(t) \cdot$$

• rot frame:
$$mP^{\dagger}(t)\frac{\mathrm{d}^2}{\mathrm{d}t^2}\left[P(t)P^{\dagger}(t)\vec{r}_{\mathrm{lab}}(t)\right] = P^{\dagger}(t)\vec{F}_{\mathrm{lab}}(t)$$

$$mP^{\dagger}(t)\frac{d}{dt}\left(P(t)P^{\dagger}(t)\frac{d}{dt}\left[P(t)\vec{r}_{rot}(t)\right]\right) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(P^{\dagger}(t)\frac{d}{dt}P(t)\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(\frac{P^{\dagger}(t)\dot{P}(t) + \frac{d}{dt}}{dt}\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(\frac{d}{dt} + \vec{\omega}(t)\times\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$\hat{z} = \hat{z}_{rot}$$

$$\hat{y}_{rot}(t)$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}$$

$$\hat{x}_{rot}(t)$$

$$\vec{r}_{lab}(t) = P(t)\vec{r}_{rot}(t)$$
$$\vec{F}_{lab}(t) = P(t)\vec{F}_{rot}(t)$$

$$P^{\dagger}(t)\frac{\mathrm{d}}{\mathrm{d}t}P(t)\vec{f}(t) = P^{\dagger}(t)\dot{P}(t)\vec{f}(t) + \frac{\mathrm{d}}{\mathrm{d}t}\vec{f}(t)$$
$$= \left(P^{\dagger}(t)\dot{P}(t) + \frac{\mathrm{d}}{\mathrm{d}t}\right)\vec{f}(t)$$

$$P^{\dagger}(t)\dot{P}(t) = \overrightarrow{\omega}(t) \times$$



Newton's law

• lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t) / P^{\dagger}(t) \cdot$$

• rot frame:
$$mP^{\dagger}(t)\frac{d^2}{dt^2}\left[P(t)P^{\dagger}(t)\vec{r}_{lab}(t)\right] = P^{\dagger}(t)\vec{F}_{lab}(t)$$

$$mP^{\dagger}(t)\frac{d}{dt}\left(P(t)P^{\dagger}(t)\frac{d}{dt}\left[P(t)\vec{r}_{rot}(t)\right]\right) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(P^{\dagger}(t)\frac{d}{dt}P(t)\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(P^{\dagger}(t)\dot{P}(t) + \frac{d}{dt}\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(\frac{d}{dt} + \vec{\omega}(t) \times\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$
$$m\left(\frac{d}{dt} + \vec{\omega}(t) \times\right)^{2}\vec{r}_{rot}(t) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

 $\hat{z} = \hat{z}_{rot}$ $\hat{y}_{rot}(t)$ \hat{x} \hat{x} \hat{x} $\hat{x}_{rot}(t)$

$$\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$$
$$\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$$

$$P^{\dagger}(t)\frac{\mathrm{d}}{\mathrm{d}t}P(t)\vec{f}(t) = P^{\dagger}(t)\dot{P}(t)\vec{f}(t) + \frac{\mathrm{d}}{\mathrm{d}t}\vec{f}(t)$$
$$= \left(P^{\dagger}(t)\dot{P}(t) + \frac{\mathrm{d}}{\mathrm{d}t}\right)\vec{f}(t)$$

$$P^{\dagger}(t)\dot{P}(t) = \overrightarrow{\omega}(t) \times$$



- Newton's law
 - ► lab frame: $m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t)$
 - ► rot frame:

$$m\left(\frac{\mathrm{d}}{\mathrm{d}t} + \vec{\omega}(t) \times\right) \left(\dot{\vec{r}}_{\mathrm{rot}}(t) + \vec{\omega}(t) \times \vec{r}_{\mathrm{rot}}(t)\right) = P^{\dagger}(t) \overrightarrow{F}_{\mathrm{lab}}(t)$$





- Newton's law
 - ► lab frame: $m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t)$

$$\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$$
$$\vec{F}_{\text{lab}}(t) = P(t)\vec{F}_{\text{rot}}(t)$$

• rot frame:
$$m\left(\frac{\mathrm{d}}{\mathrm{d}t} + \vec{\omega}(t) \times\right) \left(\dot{\vec{r}}_{\mathrm{rot}}(t) + \vec{\omega}(t) \times \vec{r}_{\mathrm{rot}}(t)\right) = P^{\dagger}(t) \overrightarrow{F}_{\mathrm{lab}}(t)$$

HW:
$$m\ddot{\vec{r}}_{rot}(t) + m\dot{\vec{\omega}}(t) \times \vec{r}_{rot}(t) + 2m\vec{\omega}(t) \times \dot{\vec{r}}_{rot}(t) + m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{rot}(t)\right) = P^{\dagger}(t)\vec{F}_{lab}(t)$$

 $\hat{z} = \hat{z}_{\rm rot}$

 \hat{x}



- Newton's law
 - ► lab frame: $m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t)$

$$\overrightarrow{F}_{\text{lab}}(t) = P(t) \overrightarrow{F}_{\text{rot}}(t)$$

 $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$

► rot frame:
$$m\left(\frac{\mathrm{d}}{\mathrm{d}t} + \vec{\omega}(t) \times \right) \left(\dot{\vec{r}}_{\mathrm{rot}}(t) + \vec{\omega}(t) \times \vec{r}_{\mathrm{rot}}(t)\right) = P^{\dagger}(t) \vec{F}_{\mathrm{lab}}(t)$$

$$\vec{m}\vec{r}_{\rm rot}(t) + \vec{m}\vec{\omega}(t) \times \vec{r}_{\rm rot}(t) + 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) + m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right) = P^{\dagger}(t)\vec{F}_{\rm lab}(t)$$

$$m\ddot{\vec{r}}_{\rm rot}(t) = P^{\dagger}(t)\vec{F}_{\rm lab}(t) - m\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) - 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) - m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right)$$

transformed original force

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 $\hat{z} = \hat{z}_{\rm rot}$

x̂ ►



- Newton's law
 - lab frame: $m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t)$

$$\vec{r}_{\rm lab}(t) = P(t)\vec{r}_{\rm rot}(t)$$
$$\vec{F}_{\rm lab}(t) = P(t)\vec{F}_{\rm rot}(t)$$

• rot frame:
$$m\left(\frac{\mathrm{d}}{\mathrm{d}t} + \vec{\omega}(t) \times \right) \left(\dot{\vec{r}}_{\mathrm{rot}}(t) + \vec{\omega}(t) \times \vec{r}_{\mathrm{rot}}(t)\right) = P^{\dagger}(t)\vec{F}_{\mathrm{lab}}(t)$$

$$m\ddot{\vec{r}}_{\rm rot}(t) + m\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) + 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) + m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right) = P^{\dagger}(t)\vec{F}_{\rm lab}(t)$$

$$m\ddot{\vec{r}}_{\rm rot}(t) = P^{\dagger}(t)\vec{F}_{\rm lab}(t) - m\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) - 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) - m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right)$$

transformed Euler original force force

 $\hat{z} = \hat{z}_{\rm rot}$

x



- Newton's law
 - ► lab frame: $m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t)$

$$\overrightarrow{F}_{\text{lab}}(t) = P(t) \overrightarrow{F}_{\text{rot}}(t)$$

 $\vec{r}_{\text{lab}}(t) = P(t)\vec{r}_{\text{rot}}(t)$

► rot frame:
$$m\left(\frac{\mathrm{d}}{\mathrm{d}t} + \vec{\omega}(t) \times \right) \left(\dot{\vec{r}}_{\mathrm{rot}}(t) + \vec{\omega}(t) \times \vec{r}_{\mathrm{rot}}(t)\right) = P^{\dagger}(t) \vec{F}_{\mathrm{lab}}(t)$$

$$m\ddot{\vec{r}}_{\rm rot}(t) + m\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) + 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) + m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right) = P^{\dagger}(t)\vec{F}_{\rm lab}(t)$$

$$m\ddot{\vec{r}}_{\rm rot}(t) = P^{\dagger}(t)\vec{F}_{\rm lab}(t) - m\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) - 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) - m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right)$$

transformed	Euler	Coriolis
original force	force	force

 $\hat{z} = \hat{z}_{\rm rot}$

x̂ ⊾



Newton's law



 $\hat{z} = \hat{z}_{\rm rot}$

▶ lab frame:
$$m \frac{d^2}{dt^2} \vec{r}_{lab}(t) = \vec{F}_{lab}(t)$$
▶ rot frame: $m \left(\frac{d}{dt} + \vec{\omega}(t) \times \right) \left(\dot{\vec{r}}_{rot}(t) + \vec{\omega}(t) \times \vec{r}_{rot}(t)\right) = P^{\dagger}(t) \vec{F}_{lab}(t)$
 $m \vec{\vec{r}}_{rot}(t) + m \vec{\omega}(t) \times \vec{r}_{rot}(t) + 2m \vec{\omega}(t) \times \vec{r}_{rot}(t) + m \vec{\omega}(t) \times (\vec{\omega}(t) \times \vec{r}_{rot}(t)) = P^{\dagger}(t) \vec{F}_{lab}(t)$
 $m \vec{\vec{r}}_{rot}(t) = P^{\dagger}(t) \vec{F}_{lab}(t) - m \vec{\omega}(t) \times \vec{r}_{rot}(t) - 2m \vec{\omega}(t) \times \dot{\vec{r}}_{rot}(t) - m \vec{\omega}(t) \times (\vec{\omega}(t) \times \vec{r}_{rot}(t))$

transformed	Euler	Coriolis	centrifugal	
original force	force	force	force	



Newton's law



 $\hat{z} = \hat{z}_{\rm rot}$

▶ ;	ab frame:	$m\frac{\mathrm{d}^2}{\mathrm{d}t^2}\vec{r}_{\mathrm{lab}}(t) = \vec{h}$	$\vec{F}_{lab}(t)$	$F_{\text{lab}}(t) = P(t)I$	$\overline{F}_{rot}(t)$
► r	ot frame:		$m\left(\frac{\mathrm{d}}{\mathrm{d}t} + \overrightarrow{\omega}(t) \times\right)$	$\int \left(\dot{\vec{r}}_{\rm rot}(t) + \vec{\omega}(t) \times \vec{r}_{\rm rot} \right)$	$_{\rm tt}(t) = P^{\dagger}(t) \overrightarrow{F}_{\rm lab}(t)$
ŀ	$n\ddot{\vec{r}}_{\rm rot}(t) + d$	$m\dot{\overrightarrow{\omega}}(t) \times \overrightarrow{r}_{\rm rot}(t) +$	$2m\vec{\omega}(t) \times \dot{\vec{r}}_{rot}(t)$ -	+ $m\vec{\omega}(t) \times (\vec{\omega}(t) \times \vec{r}_{rot})$	$_{\rm ot}(t) = P^{\dagger}(t) \overrightarrow{F}_{\rm lab}(t)$
$m\ddot{\vec{r}}_{\rm rot}(t) = P^{\dagger}(t)\vec{F}_{\rm lab}(t) - m\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) - 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) - m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right)$					
		transformed original force	Euler force	Coriolis force	centrifugal force

• fictitious forces arise from Galilean term $P^{\dagger}(t)\dot{P}(t)$

$$P^{\dagger}(t)\dot{P}(t) = \overrightarrow{\omega}(t) \times$$



• Newton's law



Þ	lab frame: 1	$m\frac{\mathrm{d}^2}{\mathrm{d}t^2}\vec{r}_{\mathrm{lab}}(t) = \overline{F}$	$\vec{t}_{lab}(t)$	$\dot{F}_{lab}(t) = P(t)F$	$r_{\rm rot}(t)$	
Þ	rot frame:	ľ	$n\left(\frac{\mathrm{d}}{\mathrm{d}t} + \overrightarrow{\omega}(t) \times\right)$	$\int \left(\dot{\vec{r}}_{\rm rot}(t) + \vec{\omega}(t) \times \vec{r}_{\rm rot} \right)$	(t) = $P^{\dagger}(t) \overrightarrow{F}_{lab}(t)$	
	$m\ddot{\vec{r}}_{\rm rot}(t) + m\dot{\vec{r}}$	$\dot{\vec{\omega}}(t) \times \vec{r}_{\rm rot}(t) + 2$	$2m\vec{\omega}(t) \times \dot{\vec{r}}_{rot}(t) +$	$- m \overrightarrow{\omega}(t) \times \left(\overrightarrow{\omega}(t) \times \overrightarrow{r}_{ro} \right)$	$_{\rm t}(t) = P^{\dagger}(t) \overrightarrow{F}_{\rm lab}(t)$	
	$\vec{m}\vec{r}_{\rm rot}(t) = P^{\dagger}(t)\vec{F}_{\rm lab}(t) - \vec{m}\vec{\omega}(t) \times \vec{r}_{\rm rot}(t) - 2m\vec{\omega}(t) \times \dot{\vec{r}}_{\rm rot}(t) - m\vec{\omega}(t) \times \left(\vec{\omega}(t) \times \vec{r}_{\rm rot}(t)\right)$					
	tra ori	ansformed ginal force	Euler force	Coriolis force	centrifugal force	

• fictitious forces arise from Galilean term $P^{\dagger}(t)\dot{P}(t)$

 $P^{\dagger}(t)\dot{P}(t) = \overrightarrow{\omega}(t) \times$

Q: can we understand dynamical stabilization as a fictitious force in some rotating frame? Marín Bukov MPI-PKS







- Rotating reference frames
 - quantum systems



lab frame

Outline

rotating frame

Quantum mechanics



- rotating reference frame
 - not inertial
 - fictitious forces
- transformation between lab and rotating frames



lab frame

rot frame

P(t)

Quantum mechanics



- rotating reference frame
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$$|\psi(t)\rangle_{\text{lab}} = P(t) |\psi(t)\rangle_{\text{rot}}$$

Quantum mechanics



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 $|\psi(t)\rangle_{\text{lab}} = P(t) |\psi(t)\rangle_{\text{rot}}$ $U_{\text{lab}}(t,0) = P(t)U_{\text{rot}}(t,0)P^{\dagger}(0)$
$\hat{z} = \hat{z}_{rot}$ $\hat{y}_{rot}(t)$ \hat{x} \hat{x} \hat{x} $\hat{x}_{rot}(t)$

- Schrödinger's equation (set $\hbar = 1$)
 - lab frame: $i\partial_t |\psi(t)\rangle_{\text{lab}} = H_{\text{lab}} |\psi(t)\rangle_{\text{lab}} / P^{\dagger}(t) \cdot$

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energy in the rot frame not the same as the transformed lab-frame Hamiltonian

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- in general, cannot evaluate in closed form
- do not have meaningful stationary states (in general)









- Periodically driven quantum systems
 - Floquet theorem
 - Floquet engineering





• time dependence H(t) = H(t + T), $T = 2\pi/\omega$

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micromotion Floquet Hamiltonian

 $-\pi$

 π

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mmm

 $-\pi$

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cf. static Hamiltonians:

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P(t) = P(t + T) micromotion Floquet Hamiltonian periodic with same period *T* as drive

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- why useful?
 - theory similar to static systems
 - time-scale separation in high-frequency limit



 $-\pi$

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• physical meaning of Floquet's theorem? $U(t,0) = P(t) \exp(-itH_F)$

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► recall: $U_{\text{lab}}(t,0) = P(t)U_{\text{rot}}(t,0)P^{\dagger}(0)$



lab frame

rot frame

P(t)





defined by the micromotion operator P(t)

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lab frame

rotating frame



 $H_{\rm rot}(t) = P^{\dagger}(t)H_{\rm lab}P(t) - iP^{\dagger}(t)\dot{P}(t)$

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lab frame

H(t)

 $P^{\dagger}(t)$

rotating frame

 $H_{\rm F}$



- Floquet's theorem proves the existence of a special rotating frame, defined by the micromotion operator P(t)
 - the rotating frame Hamiltonian is the time-independent ${\cal H}_{F}$

$$H_F = P^{\dagger}(t)H(t)P(t) - iP^{\dagger}(t)\dot{P}(t)$$

• note: H_F contains fictitious force potential $iP^{\dagger}\dot{P}$!



• Floquet engineering: how do we choose the drive H(t) to design properties to H_F ?

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- Examples
 - spin-1 particle in a circularly polarized drive
 - quantum Kapitza oscillator
 - artificial gauge fields



Spin-1 particle in a circularly polarized magnetic field



► spin-1 matrices

$$S^{x} = \frac{1}{\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} \qquad \qquad S^{y} = \frac{1}{i\sqrt{2}} \begin{pmatrix} 0 & 1 & 0 \\ -1 & 0 & 1 \\ 0 & -1 & 0 \end{pmatrix} \qquad \qquad S^{z} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & -1 \end{pmatrix}$$



Spin-1 particle in a circularly polarized magnetic field

- Hamiltonian $H(t) = \Delta S^z + g(\cos \omega t S^x + \sin \omega t S^y)$
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 $\langle \vec{\sigma} \rangle$

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$$P(t+T) = e^{-i\omega(t+T)S^{z}} = e^{-i\omega t S^{z}} e^{-i\omega T S^{z}} = e^{-i\omega t S^{z}} = P(t)$$

 $P^{\dagger}(t)S^{z}P(t) = S^{z}$

► can show (HW): $P^{\dagger}(t)(\cos \omega t S^{x} + \sin \omega t S^{y})P(t) = S^{x}$ (by design) $iP^{\dagger}(t)\dot{P}(t) = \omega S^{z}$



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spin-1 matrices

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• define co-rotating frame $P(t) = \exp(-i\omega t S^z)$

► check:
$$P(t+T) = e^{-i\omega(t+T)S^{z}} = e^{-i\omega t S^{z}} e^{-i\omega T S^{z}} = e^{-i\omega t S^{z}} = P(t)$$

 $P^{\dagger}(t)S^{z}P(t) = S^{z}$

- can show (HW): $P^{\dagger}(t)(\cos \omega t S^x + \sin \omega t S^y)P(t) = S^x$ (by design) $iP^{\dagger}(t)\dot{P}(t) = \omega S^z$
- Floquet Hamiltonian: $H_F = P^{\dagger}(t)H(t)P(t) iP^{\dagger}(t)\dot{P}(t) = (\Delta \omega)S^z + gS^x$

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Spin-1 particle in a circularly polarized magnetic field



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- note: very few exactly solvable models












- Examples
 - quantum Kapitza oscillator
 - artificial gauge fields



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 - stable inverted equilibrium for $A \gg \omega_0$



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$$H_{\rm rot}(t) = \mathscr{P}^{\dagger}(t) \left(\frac{p^2}{2} - \left(\omega_0^2 + \underline{A\omega \cos \omega t} \right) \cos \theta \right) \mathscr{P}(t) - \underline{i} \mathscr{P}^{\dagger}(t) \dot{\mathscr{P}}(t)$$



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HW:
$$\mathscr{P}^{\dagger}(t)p\mathscr{P}(t) = p - i\partial_{\theta}\left(iA\sin\omega t \cos\theta\right) = p - A\sin\omega t \sin\theta$$

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compute period-average
 effective period

effective potential $V_{\rm eff}(\theta)$

$$H_F^{(0)} = \frac{1}{T} \int_0^T H_{\text{rot}}(t) \, \mathrm{d}t = \frac{p^2}{2} + \frac{A^2}{4} \, \sin^2\theta - \omega_0^2 \, \cos\theta$$

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leading-order effective Hamiltonian

inverse-frequency expansions

$$H_F^{(0)} = \frac{1}{T} \int_0^T H_{\text{rot}}(t) \, \mathrm{d}t = \frac{p^2}{2} + \frac{A^2}{4} \sin^2 \theta - \omega_0^2 \, \cos \theta \qquad \qquad H_F = \sum_{n=0}^\infty \omega^{-n} H_F^{(n)}$$

effective potential $V_{\text{eff}}(\theta)$
$$P(t) = \prod_{n=0}^\infty P^{(n)}(t)$$

 $P^{(0)}(t) = \mathscr{P}(t)$

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• Hamiltonian
$$H(t) = \frac{p^2}{2} - (\omega_0^2 + A\omega \cos \omega t) \cos \theta$$

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• leading-order effective Hamiltonian

$$H_{F}^{(0)} = \frac{1}{T} \int_{0}^{T} H_{rot}(t) dt = \frac{p^{2}}{2} + \frac{A^{2}}{4} \sin^{2}\theta - \omega_{0}^{2} \cos\theta$$
effective potential $V_{eff}(\theta)$
• analyze stability at $\theta = \pm \pi$
 $\partial_{\theta}^{2} V_{eff}(\theta) = \omega_{0}^{2} \cos\theta + \frac{A^{2}}{2} \cos 2\theta$
 $\int_{\theta} = \pi$
Marin Bukov
• leading-order effective Hamiltonian
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Outline



• Examples

artificial gauge fields







• compare:
$$H_{\text{rot}}(t) = \frac{1}{2} \left(p - A \sin \omega t \sin \theta \right)^2 - \omega_0^2 \cos \theta$$
 vs

Kapitza pendulum in rotating frame

vs. $H = \frac{1}{2} (p - A(x))^2 + V(x)$

particle in magnetic field?

• compare: $H_{\text{rot}}(t) = \frac{1}{2} \left(p - A \sin \omega t \sin \theta \right)^2 - \omega_0^2 \cos \theta$ vs. I

$$H = \frac{1}{2} (p - A(x))^{2} + V(x)$$

gauge potential but no magnetic field in 1d!

particle in magnetic field?

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 - gauge potential but no magnetic field in 1d!
- compare: $\vec{F}_{\text{Coriolis}} = -2m\vec{\omega} \times \vec{v}$ vs. $\vec{F}_{\text{Lorentz}} = -q\vec{B} \times \vec{v}$
 - artificial magnetic fields from *Floquet engineering*

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artificial magnetic fields from Floquet engineering



superconductor: Abrikosov vortex lattice



Abrikosov, Nobel Lecture, Rev Mod Phys 76 975 (2004)

 $NbSe_2 \\ \mbox{type-II superconductor}$

scanning tunneling microscopy (STM)

VS.



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artificial magnetic fields from *Floquet engineering*



ultracold atoms in optical lattices

- quantum simulation of topological insulators
- *but:* no orbital *B*-field effects for neutral atoms

Floquet engineered magnetic fields







MAX PLANCK INSTITUTE FOR THE PHYSICS OF COMPLEX SYSTEMS

Floquet engineering for quantum simulation



MPI-PKS, Dresden

- Floquet engineering: periodic drives ascribe new properties to physical systems
 - dynamical stabilization
 - artificial gauge fields (topological insulators, etc.)
- caveat: driven systems may absorb energy (heat death)

key idea: design fictitious forces in rotating reference frame



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THE PHYSICS OF COMPLEX SYSTEMS

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